

1 **On the Time Decay constant of AEM systems: A Semi- Heuristic Algorithm to validate**
2 **calculations.**

3
4 José Manuel Martínez
5 2-572 Ampere Ave.
6 Laval, Canada, H7N5W9
7 Phone: 438-22-3366
8 E-mail: jmartinez2050@gmail.com

9
10 Richard Smith
11 Harquail School of Earth Sciences
12 Laurentian University
13 935 Ramsey Lake Road, Sudbury, ON, Canada, P3E 2C6, RSSmith@laurentian.ca
14 Danilo Diaz Vazquez
15 Departamento de Ciencias Fisicas, Universidad Andres Bello, Autopista Concepcion-
16 Talcahuano 7100, Talcahuano, Chile
17 Email: danielodiaz@unab.cl

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31 **ABSTRACT**

32 The time decay constant or “tau” of airborne EM systems (AEM) is commonly used to indicate
33 the presence and the relative conductivity or conductance of conductors in the survey area. In
34 fact, it is not a constant since it depends on: the system; the survey design and the method of
35 calculation. The system dependence is a consequence of parameters relating to the acquisition
36 and *pre-* and *post-* processing of the signal. Here we propose a method for calculating tau,
37 which is simply the time at which the transient voltage decays to 37%, or V_{37} , of some initial
38 value. The model utilises a semi-heuristic algorithm that estimates V_{37} for each transient in the
39 database and then it calculates the delay time that voltage is measures, which is the estimates
40 tau value. No calculation is involved with the data, instead tau is given by a weighted average
41 of the delay times associated with the windows either side of the V_{37} value. We illustrate how
42 this algorithm works using data collected by MEGATEM II in the Reid-Mahaffy test site.
43 Results shown a good agreement between tau-grids reported in previous studies and those
44 calculated with our V_{37} – method. To account for all effects coming from the acquisition and
45 processing of EM data, the algorithm allows the emphasis to be shifted from early to late time
46 parts of the transient. It is envisage that since this method does not apply any mathematical
47 operation to the data it may serve as a robust means of for validating other methods.

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51 **Key Words:** time decay constant, tau, semi-heuristic method, large EM database, fast
52 calculation, time domain.

53

54 INTRODUCTION

55

56 A map of tau with electromagnetics (EM) anomalies superimposed is a common derivable of
57 AEM surveys, as it summarizes the EM signature of the area. One of the first heuristic methods
58 for calculation of tau was reported by Nelson and Morris (1969) where the tau was calculated
59 using the ratio between two late off-time windows. In this case, the interpreter needs to judge
60 by experience which late time windows are suitable for calculations. Despite of problems, the
61 methodology could be used to summarize an entire survey. Grant and West (1965) devised a
62 method to determine the tau for frequency-domain data by taking the ratio between the in-
63 phase and quadrature responses. More insight into the nature of transient decays was provided
64 by Palacky and West (1973), who showed that these decay transients are described at early
65 times ($t < \tau$) by a power law while for late times ($t > \tau$) by multiple exponential decays.
66 This approach has since been extensively applied to understand and explain both synthetic and
67 field EM data. Further studies carried out by Kaufman (1965, 1978) demonstrated that for late
68 times ($t > \tau$) a single exponential function was appropriate to describe the presence of deep
69 seated conductive bodies. Consequently, the single exponential and power-law behaviour of
70 the decay gained popularity and it was further studied by Bartel and Hohmann (1985) and
71 Bartel and Becker (1988), and others. These studies undertook forward modelling to establish
72 very useful empirical formulas for calculating tau based on: the terrain clearance, depth, size
73 and conductivity contrast. Other methods like singular value decomposition into multiple
74 exponentials by Stolz and Macnae, (1998) and the matrix pencil method by Chen and Macnae
75 (1998) are good trade-off between time of computation and the reproducibility of τ . Currently,
76 it is widely accepted that late transients are well described by a single exponential decay.
77 Commonly, tau grids are calculated by fitting a single exponential decay curve to the EM data,
78 e.g. Smith and Lee, (2001).

79 The aim of this study was to find a method to determine tau with as few operations to the data
80 as possible. To that end we investigate how EM data is acquired from current AEM systems
81 and found out that there are several engineering trade-offs made during the making of an AEM
82 system that impact on the decay transients and thus tau. Upon analysing these effects our semi-
83 heuristic method to calculate tau for large EM datasets emerged. The method is based on the
84 late-time transients being described by a single exponential function. In that case tau is simply
85 the time at which the voltage measured by receivers decay to 37% of its initial value. Therefore,
86 an algorithm seeks this value, denoted V_{37} , for each transient in the database. The tau value is
87 a weighted average of the delay times of the two windows between which the value of V_{37}
88 occurs (described in more detail below). There is not mathematical operation involving the EM
89 measurements, so measurement noise does not have a large impact on the estimated tau value.
90 The worse impact usually no worse than moving the tau estimate one time window earlier or
91 later.

92 **The rationale behind tau measurements**

93 Why is tau EM system dependent?

94 The tau is system dependent because of both the transmission and reception of primary and
95 secondary fields are system dependent. To illustrate this strong dependence, let's consider the
96 two conditions for which tau is system independent. They are: 1) a perfect square pulse (i.e.
97 rise and fall times occur instantaneously) with at least a 50% duty cycle and 2) voltages
98 measured by receivers that are only measuring the voltages in the off-time when there is no
99 primary field. If these two ideal conditions are met, then the decay transients are those
100 illustrated in Figure 1B.

101 {INSERT FIGURE 1 HERE}

102 In the idealized waveform of Figure 1, where the current rises from zero to I_{peak} instantaneously
103 and falls instantaneously at the end of the pulse. There is a transient response from the
104 target for every time a current switch occurs. Note that in the entire full-cycle, T there is four
105 responses (i.e. one occurring after each switching). These responses are not hypothetical they
106 are present in all AEM systems but they are not always recorded, for reasons discussed later.
107 In our illustrative example on Figure 1B, the same target is excited, so there is no differences
108 in these decay transients, i.e., $\tau_1 = \tau_2 = \tau_3 = \tau_4$. When recording over a whole duty cycle, T ,
109 the tau can be determined by stacking these four responses and apply any of the previous
110 methods described in the introduction. Or alternatively, one can estimate the maximum voltage
111 response from the target, i.e. V_0 and use this to calculate the V_{37} , which is 37% of this initial
112 value and determine the time that this occurs, a process which will be described in more detail
113 later on. Note that only and only if there is perfect reception where the voltage measured by
114 the receiver are solely due to the responses from targets, will both the decay transients and thus
115 tau be system independent. In practice it is not possible to switch the current on or off
116 instantaneously. Hence, the secondary will not be the only field present in the measurements,
117 there will also be a primary field. The reason for the non-ideal nature of waveforms is that the
118 current has to pass through cables, and they have inductances, which introduce a delay to the
119 switching process. This means that there is a changing primary field, which will induce a
120 voltage in the receiver coil. This primary field can six orders of magnitude larger than the
121 voltages from the response of targets. This is attempted to be addressed by introducing an
122 attenuation and/or bucking mechanism so as to literally avoid burning the conditioning
123 electronics of the receiver. Each electronic engineer designing EM systems has their own
124 approaches of acquiring EM data to avoid these problems and specific approaches to pre and
125 post processing data. These approaches can make tau system dependent. Hence one may
126 wonder what percentage of a tau estimate reflects the target and not the system. In our view,

127 the more clarity we gain about how all these factors interplaying, the better tau can be
128 associated with the target and exclude the system and/or method of calculation.

129 In practice, the waveform of an AEM system is driven by some form of square pulse as depicted
130 in Figure 2.

131 Effect of synchronous acquisition

132 The cornerstone of acquisition of existing AEM systems relies on synchronous acquisition
133 schemes which utilises gates (Macnae et al., 1984), such as those in Figure 2 labelled W1, W2,
134 W3, and so on. Throughout these windows the analog signal picked by the receiver sensors can
135 be converted into digital form (details are given in the next section). One can imagine when
136 these gates are open during the off-time, the data acquisition system (DAS) collects samples,
137 N_s (see open circle in Figure 2B) depending on the system sampling rate, SR (more details of
138 which are given later on). For instance, if the DAS has an analog module rated at $SR = 250$ kHz,
139 the very first sample gathered by the system will come after $4\mu s$ (this calculation is performed
140 below) after the primary field has been shut down. These windows are designed such that the
141 number of samples in each window increases exponentially for later windows W_i . It is
142 expected that open circles in each window are averaged in some way i.e., $\sum_{i=1}^N V_i / N$. The simple
143 average works well if the window is narrow and the response decays approximately linearly in
144 the window. However, for late time ($t > \tau$) some other statistical estimator, like the median
145 may offer a better statistical descriptor (p.656 Press et al., 2007). However, there might be
146 other undocumented processes used to improve the results like using more samples from
147 outside the windows, or arbitrarily removing a late-time asymptote or a non-decaying
148 component (as is done in the TEMPEST system; Lane et al., 2000). Obviously, the procedures
149 used to process the raw data will determine the decay transients and thus tau.

150 {INSERT FIGURE 2 HERE}

151 Effect of receiver coils

152 Transients at the receiver (Rx) falloff rapidly lasting just a few milliseconds before they sink
153 into noise. Electrically Rx are designed to have low noise figures (in the order of pT) and flat
154 responses with a bandwidth of about 20kHz (Macnae, 2012). The most desirable electrical
155 characteristics of an Rx is to have low impedance (in order to have low noise figures) and the
156 highest inductance possible (to be able to pick small varying secondary magnetic fields). These
157 two physical properties are contradictory, leading to a trade-off between inductance and
158 impedance for an inductive receiver at low frequency ranges (Martinez et al., 2014). High
159 inductance means smaller wire gauges and lots of turns and that increase the impedance of the
160 coil drastically. One solution to attain high inductances in the Rx, with fewer turns, is to
161 utilising permeable core materials to amplify the magnetic field. The only inconvenience is that
162 these materials impart non-linearity to the pre-amplification stages (Neri, 1996). A close-loop
163 compensation network is necessary to flatten the electromagnetic field responses within the
164 bandwidth. Another solution, since the peak voltage from targets, i.e., V_o depends on r_{RX}^2 ,
165 would be to increase the Rx radius, r_{RX} . However, since the voltage induced by the primary
166 field, V_p is up to six orders of magnitude larger than V_o this radius is restricted to $r_{RX} \sim 0.6m$.
167 Note that, V_o is determined by the physical dimension of the Rx and it is nothing but $2\pi\mu_0 A_{eff}$,
168 where μ_0 and A_{eff} are the permeability of the free space and the effective area, $A_{eff} =$
169 $\pi r_{RX}^2 N_{RX}$ of the Rx, respectively. Because of these differences between V_p and V_o a dual
170 preamplifier is needed in an AEM systems. The amplifier for the secondary field attenuates
171 during T_{ON} while amplifies during T_{OFF} . This strong attenuation impedes recording of the on-
172 time response from the target

173 atone limitation of the receiver is that there is a lower limit of voltage resolution, V_R for each
174 AEM system. This is, the minimum voltage that can be converted from analog to digital. In
175 other words, V_R depends on the ADC bit resolution of the DAS. For instance, if the DAS has

176 an analog module capable of measuring voltages between 0 and 5V and the ADC encodes its
 177 data using 16 bits, this will give: $V_R = 5V/2^{16} = 76.3\mu V$. In other words the lowest voltage the
 178 EM system can detect is just $76.3\mu V$ and the maximum is less than 5V. Hence several pre-
 179 amplification stages are required. Knowing the total gain, G of the EM system would be
 180 sufficient to obtain transients reported in the EM database from the raw data. In some cases the
 181 measurements are normalized, with respect to primary magnetic fields recordings, e.g.
 182 $B(\text{ppm}) = V_o G / V_p \times 10^6$ in [nT/V]. Sometimes V_p may be expressed in ADC units
 183 ($V_p(\text{ADC}) = 2 \times 10^4 \text{ADC}$), that one can convert to voltage using V_R . That is: $V_p(\text{V}) = 2 \times$
 184 $10^4 \text{ADC} \times 76.3\mu V / \text{ADC} = 1.5\text{V}$.

185 The above briefly describes the choices made for how an AEM system acquires data, and these
 186 can impact how decay transients are recorded.

187 Effect of data encodings

188 One way to improve the SNR is by acquiring more samples N_s per windows sampled (Eaton,
 189 2013) by increasing the system's sampling rate, SR . For instance, if the DAS of the EM system
 190 is rated at 40 MHz with one tick every 25ns; having an analog module rated at $SR = 250$ kHz;
 191 its first sample will come after $4\mu s$ ($= \frac{40\text{MHz}}{0.25\text{MHz}} \times 0.025\mu s$). The latter is some time referred to
 192 as the as point value. Nowadays, the SR of commercial and non-commercial system is set
 193 around $SR = 250$ kHz (Allard, 2007). Ramping up the SR up is not a straight-forward task as
 194 the volume of data to be streamed is in the order of megabit/s and special chopping and
 195 streaming algorithms are needed right on the memory of the microcontrollers to cope with such
 196 a huge volume of data (Eaton, 2013). To have an idea of the order of streaming happening in
 197 an AEM system, if the system has a $SR = 500$ kHz and the waveform repeats at a base frequency,
 198 $f_b = 30$ Hz; then $N_s = \frac{SR}{2f_b} = \frac{500\text{kHz}}{60} = 8333.3$ samples per $T_{T/2}$. Let's say these samples are
 199 encoded using 32 bits, then the number of bits in $T_{T/2}$ is 266665.6 (8333.3×32) bits. As there

200 are 60 half cycles per second, then in 1s there are 1,599,993.6 (266665.6x60) bits (or 2Mb/s)
201 and if the system has three Rx then it will be 6MB/s. That may not sound much data in today's
202 line communication systems, however the catch for an AEM system is that the DAS handles
203 several mixed signals (analog & digital signals) requiring algorithms to ensure precise
204 triggering, synchronization, streaming and control. And all of them have to reside on the
205 memory of the microcontroller to handle all system request. Therefore control, communication
206 and steaming protocols needs to be optimized in order to avoid overflow, lags so as to ensure
207 than data is not lost.

208 All the above bring some clarity on the way the AEM system acquire convert and encode its
209 data to form the decay transients. Manufacturer are quite familiar with these issues but not all
210 interpreter are. However, it is important to know these issues in order to understand how tau
211 depends on a particular AEM system and how to compare the tau gathered from two different
212 AEM systems.

213 Effect of the transmitter moment + switching

214 Significant engineering efforts are required to increase the moment (power) of an AEM, where
215 the moment, $m = NIA$. This is as important because: 1) theoretically the measured signal
216 depends linearly with m and 2) in practice by doubling the duty cycle of the current wave form
217 will guarantee higher target responses (Becker et al., 1984). Often there is a tendency to use
218 the peak current, I_{peak} , without including the duty cycle, in order to calculate m (see for
219 instance I_{peak} and I in Figure 2B). The reason of this widespread confusion comes from the
220 fact that in frequency domain systems the current is always ON, i.e. 100% duty cycle. For AEM
221 system this is not the case. For instance, for the MegaTEM system, $I_{peak}=710$ A and when the
222 system is tuned to fire at 30 Hz and 90 Hz base frequencies, their corresponding duty cycle are
223 24% and 41%, respectively. Details regarding this calculation is provided in the result section.

224 This means the current for these frequencies are $I = 710 \times 0.24 \text{ A} = 170 \text{ A}$ at 30 Hz. This value
225 to be entered in the calculation of m and not $I=710 \text{ A}$.

226 The tau value can change depending on the shape of the waveform, which elicits responses
227 from targets by creating a sharp rise or fall of the pulse of current (see the ideal situation in
228 Figure 1A). The current is induced in a target according to Faraday's law, which states the
229 electric field induced by the primary is directly proportional to the rate of change of the primary
230 pulse, i.e. $\Delta I_p / \Delta t$ (see for instance, slow and fast switching in Figure 2A). The sharper $\Delta I_p / \Delta t$,
231 the more high frequency harmonics are imparted into the target. To obtain sharper rises or falls
232 (via shorter ramps) is one technical aspect that is difficult to achieve in an AEM system. This
233 is as such because: 1) the switching mechanism behind generation of primary magnetic fields
234 need to be fully digital and controlled at frequencies of at least 500 kHz (Saini et al., 2014),
235 and 2) the total inductance of the transmitter, Tx, (including transmitter cables) are rarely
236 optimized.

237 At this point one may wonder why a fast switching is of importance for tau? The answer is
238 given by comparing decay transients from slow and fast switching (see black and red decay
239 transients in Figure 2B). A fast current cut-off comes with shorter t_{RAMP} , (i.e., the time the EM
240 system takes to bring the current up or down to zero). One way to achieve a t_{RAMP} of a few
241 microseconds is by imparting few Amperes into the Tx. That is the idea behind the dual
242 moment deployed by the Skytem system described by Sørensen and Auken (2004). Obviously
243 if t_{RAMP} takes too long, then high-frequency or near-surface information is reduced. For
244 instance, in the illustrative example of Figure 2B, the first four off-time windows (W9-W11)
245 available for a fast switching system (black decay curve) are smaller amplitude when they
246 eventually occur in W12-W14 in the case of the slow switching transmission system (red decay
247 transient).

249 Effect of bucking/shielding mechanisms

250 Another difficulty with EM systems is that the voltage of transmission, V_p can be six orders of
251 magnitude greater than the secondary. The variety of solutions to this issue results in a wide
252 variety of EM systems configurations. One solution is to use a dual preamplifier. One
253 amplifier will attenuate the signal during T_{ON} when the primary is large and the second will
254 amplify it during T_{OFF} when the small primary is present. Another solution used by the VTEM
255 systems is to cancel the primary fields in a receiver by injecting the same pulse of current but
256 with reverse polarity into a more proximal transmitter bucking coil (Kuzmin and Morison,
257 2013). The disadvantage of this solution is that m will be slightly lower, compared to the
258 situation where there is no bucking coil, due to the field of the bucking coil, opposing that of
259 the main transmitter therefore lowering the overall system moment (Xiao et al. 2017). Yet
260 another solution is to null couple the primary field with the receiver, for example above the
261 transmitter wire where the field is horizontal and a receiver loop is horizontal. Of course, the
262 problem can also be avoided by placing the receiver a larger distance from the transmitter for
263 example ~40 m above and 20 m off centre.

264 **Tau & the survey design**

265 Adjusting the survey design is another way to improve the signal-to-noise ratio (SNR) of the
266 AEM system.

267 For example say T_s is the survey sampling time, which is the time elapsed between two set of
268 lateral (i.e. along the flight direction) measurements. And the distance between these
269 measurement is d . The former sometimes is referred as the lateral resolution of the survey.
270 Then, the sampling frequency of this survey is: $f_s = \frac{v_{heli}}{d}$; where v_{heli} is the speed of the
271 helicopter. The number of transients within each survey sample is $N_t = \frac{2f_b}{f_s}$. Suppose an AEM
272 survey is conducted at a base frequency $f_b = 30$ Hz; and the transients are stacked over a

273 sample interval of $T_s = 1$ s, so $f_s = 1/T_s = 1$ Hz). Then, if v_{heli} is on average 40 knots or 21m/s,
274 then will result in d of ~ 21 m and there will be $N_t = \frac{2 \times 30}{1} = 60$ transients available for stacking
275 for every 21m flown. If we double v_{heli} to 42 m/s while keeping $T_s = 1$ s; N_t will remain
276 unchanged and from the spatial resolution will deteriorate as d increases to 42 m and noise due
277 to vibration at higher speeds sink the signal into the noise way sooner. One way to improve the
278 survey while keeping v_{heli} fixed e.g., at 21 m/s is using $T_s = 4$ s). In that case, $f_s = 0.25$ Hz
279 resulting in $N_t = 240$ transients for every 42 m flown.

280 **Effect of filtering and stacking**

281 While in flight an AEM system is subject to motion, vibrations and bird swings (Davis, 2008).
282 Typically, AEM systems have an analog low-pass anti-alias filter for each receiver at the input
283 of the preamplifier to remove high-frequency components. In addition most of them remove
284 the 50Hz/60Hz power-line signal by digitally stacking the data at a frequency that has the
285 power-line signal as an even harmonic. Once the signal is in a digital mode the most common
286 workflow of processing is removing unwanted spherics and then apply a series of non-linear
287 and linear filters to the EM database. Details about the type of filters and their characteristics
288 are rarely disclosed and without doubt they impact the transients and thus tau. This workflow
289 can deteriorate the main peak of anomalies. Typically, longer duration filters reduce the high-
290 frequency noise floor noise but at the expense of decreasing the main anomaly peaks. There is
291 an alternative workflow and deals with our illustrative example of Figure 2B. It is rather
292 convenient to estimate the transient by its corresponding windows by proposing a statistic
293 descriptor for each window to bring the open circle (see Figure 2B) to full dots before the first
294 stacking occurs. In that way one can reduce, if not entirely avoided all filtering the EM data is
295 subject to. If each transient is process in that manner in every half-cycle; the next will be
296 removing its trends (see straight line draw in Figure 2B) to make stacking possible. As for
297 stacking concerned, one can use the mean value or median depending on the noise level of the

298 system. To the best of our knowledge this workflow strategy based on the actual acquisition of
299 the EM signal never has been used in the processing of the data acquired by AEM system.
300 Indeed, it is more appealing as these window filters remove the high frequency noise, but use
301 the shortest length possible.

302 **V₃₇ – METHOD**

303 Since the late time response is described by a single exponential function (Kaufman, 1978), the
304 late time voltage response for the step and impulse excitations are respectively:

305 {INSERT EQUATION 1, HERE}

306 One way to determine tau is by estimating the time at which the step response signal has
307 decayed to 37% of its initial value, V_0 . That is if t_j is the window centre time of the j th window

308 (see Figure 2B), i.e. $t_j = \frac{t_{j-1} + t_{j+1}}{2}$, then if $t_j = \tau$ for the step response yields: $V(t_j = \tau) =$

309 $V_{37} = \frac{V_0}{e} = 0.3679V_0$. Figure 3 illustrates the scheme used in this paper for computation of τ .

310 Note that on the x-axis symbols t_{j-1} and t_j represent the two consecutive centred time windows

311 where the amplitudes of the decay at these times are above and below, V_{37} , i.e. $V_{j-1} > V_{37} >$

312 V_j . Rarely does τ corresponds with a centre time that is given in the time window table. One

313 may think to calculate τ simply by taking the average time from which V_{37} falls into, i.e. $t_m =$

314 $(t_{j-1} + t_j)/2$, however for the situation depicted in Figure 3 that would be an overestimate.

315 The other option is to use a simple weighted-average interpolation. That is:

316 {INSERT EQUATION 2, HERE}

317 Weights ω_j and ω_{j-1} are values between 0 to 1 and they give a measurement of proximity of

318 V_{j-1} and V_j with regards to V_{37} . For instance, in the example of Figure 3, V_{37} falls closer to the

319 measurement V_{j-1} than V_j . That is:

320 {INSERT EQUATION 3, HERE}

321 {INSERT EQUATION 4, HERE}

322 Maximum weights are equal to 1 and this is when $V_{j-1}=V_{37}$ or $V_j =V_{37}$. In the case τ is given
323 by the central time gate and not calculated by equation (3). Otherwise these weights are less
324 than one and the higher they are the closer τ will be from either t_{j-1} or t_j (see Figure 3).

325 {INSERT FIGURE 3, HERE}

326 Exponential functions decay monotonically, so it is not physical possible that the decay voltage
327 at time “ $t + 1$ ” be greater than at time “ t ”. Therefore the first thing the V_{37} – method does is to
328 excluded any transients that do not monotonically decay.

329 **RESULTS**

330 The Reid-Mahaffy site is a test site in Canada is an area covered by glacial overburden and
331 underlain by Archean (~2.7 b.y.) mafic to intermediate metavolcanic rocks in the south, and
332 felsic to intermediate metavolcanic rocks in the north, with roughly an east-west-striking
333 stratigraphy (Ayer and Trowell, 1998).

334 The V_{37} -method was tested on two available EM datasets provided by the Ontario Geological
335 Survey (OGS 2003, 2004) collected over this site. The region was surveyed twice by the
336 MEGATEM II time-domain electromagnetic system. The 2003 data set were collected with a
337 base frequency $f_b= 90$ Hz, while the 2004 data set used $f_b=30$ Hz. In both cases, the survey
338 area was covered using sixteen N-S oriented traverse lines; with a survey line spacing of 200
339 m. The aircraft was flown at a nominal transmitter altitude of 120 m and a nominal receiver
340 altitude of 70 m. For this system the first five windows corresponds to on-time data, then there
341 are two windows in the off ramp (6 and 7) and finally 13 off-time windows. The sample rate
342 (SR) of the system was 43.4 μ s at 90 Hz and 16.3 μ s at 30 Hz. The duty cycle was not reported,

343 however it can be calculated using f_b , the on/off-time and the delay reported, see for instance
344 p.8 on. OGS (2003) for the 90 Hz case. The off and on times are 3175 μs and 2280 μs ,
345 respectively and the anti-alias filter delay was 100 μs and this amount should be added to the
346 off-time Hence the half-cycle is $2280+100+3175 = 5555 \mu\text{s}$, and therefore the duty cycle is
347 41% [$2280/5555$]. The peak current was reported as 710 A, therefore the nominal current is I
348 $= 291 \text{ A}$ ($710 \times 0.41 \text{ A}$). The transmitter has 5 turns and the area of receiver coils was 406 m^2 .
349 Therefore, the duty-cycle corrected moment for this survey was $m=591\ 600$ ($120 \times 5 \times 406$)
350 Am^2 . Similar calculation can be performed for $f_b=30 \text{ Hz}$, where the on time is typically 4000
351 μs and the half-cycle 16666 μs yielding a duty cycle of 24%; with the dipole moment 927 000
352 Am^2 , the duty-cycle corrected dipole moment is $m = 222\ 000 \text{Am}^2$. Other parameters of the
353 survey can be calculated using the formulations provided in earlier sections and their value are
354 depicted in the Table below.

355 {INSERT TABLE, HERE}

356 In the report (OGS, 2004), the EM database “darm1” (at 30 Hz) had several windows for tau
357 calculated. These tau reported were calculated from W8 to W20, which corresponds to 369 μs
358 to 3038 μs after turn-off of the primary field. The standard method used is to fit an exponential
359 $y_{\text{data}}=y_0e^{(-t/\tau)}$ to the EM decay [see p. 13 OGS (2003)].

360 The comparison shown thereafter deal with the tau reported in the OGS (2003 and 2004)
361 database window and the one calculated using equation (2). To start the comparisons, the
362 channels of the vertical component magnetic-field B_z (em_bz_final_off) was selected. In the
363 study the earliest channel parsed to the V_{37} -code was window 8. The tau derived was gridded
364 using minimum curvature, which available in many EM database processing software
365 packages.

366 {INSERT FIGURE 4 HERE}

367 The only user adjustable parameter is a tolerance to control the number of readings greater than
368 the previous reading. On figure 5, the tolerance was set to 0%, so any transient what showed
369 an increase in any window compared with the previous window was excluded from the tau
370 calculation and the tau value is set to zero.

371 {INSERT FIGURE 5, HERE}

372 Clearly, the data from the region just to the north east of the centre has zero tau and an
373 inspection of the database shows that for this area the values of window W20 are higher than
374 W19. By adjusting the tolerance to 15% none of the transients were penalized.

375 Another feature of the algorithm is that the interpreter can choose which window to assign the
376 value of V_0 before the calculation of tau. To illustrate the impact of different initial windows,
377 we selected the Bz_channel from EM database “darm2”, which is the 90 Hz) database from
378 OGS (2004) and calculated tau using window 8 and then window 9 as the initial window Figure
379 6 and Figure 7 are the grid of the tau calculated using equation (2). In this dataset, the tolerance
380 permitted was set at 5%. Note that the only effect of this tolerance is to include more transients
381 into the calculation of tau and it does not alter the estimate in any form.

382 {INSERT FIGURE 6 HERE}

383 {INSERT FIGURE 7 HERE}

384 Next we apply the method to the dB/dt data set.

385 Figure 8 illustrates the tau grid of the channel of “em_z_final_off” reported in the darm2
386 database, OGS(2004) while Figure 9 the one calculated with a tolerance of 5% using the
387 equation (2).

388 Looking at tau grids calculated (see Figures 5,6,7 and 9) we observe that there are numerous
389 central lows with peaks to the north and south. This is a consequence of the anomaly over a
390 vertical conductor having a peak-trough-peak character on the z component (Smith and

391 Keating, 1996). Hence we also tried the new method on the x component, which has a large
392 positive peak directly over a vertical conductor (Smith and Keating, 1996).

393 Figure 10 shows the standard tau grid of the channel “em_x_final_off” delivered in the darm2
394 database (OGS 2004) while Figure 11 is the tau grid calculated using (2) with an overall
395 tolerance of 5%. Figure 11 shows less along-line asymmetry artefacts and more features
396 running along the geological strike WNW.

397

398 {INSERT FIGURE 10 HERE}

399

400 {INSERT FIGURE 11 HERE}

401

402 **DISCUSSION**

403 The tau grids calculated using the standard exponential fitting (e.g. Figure 4) and those
404 calculated by the V_{37} -method (see Figure 6) are in agreement. Both show maximum at the same
405 position. There are slight differences in the value of the time constant, possibly as a
406 consequence of the standard method fitting noisy data at late time or the fitting emphasizing
407 the large (early-time) amplitudes.

408 Being able to select the first window to start and the tolerance applied to transients provides
409 interpreters with a flexible way of calculating tau so as to enhance features of interest. The entire
410 calculation takes fractions of seconds for a large EM dataset, therefore different scenarios can
411 be tested quickly. For instance, Figure 6 illustrates the tau calculated using the start window 8
412 (W8) while Figure 7 uses W9. For example, Figure 7 does a better job highlighting anomalies
413 appearing on the north-east side of the map which could be deeper. Same applies to these
414 anomalies found in the central-east side since they EM signatures even when the window to
415 start was W9.

416 In processing the above, we found that different tolerances were required for darm1 (15%)
417 compared to darm2 (5%). This is because the data in darm1 was acquired with MegaTEM
418 operating at a base frequency of 30 Hz compared to 90 Hz for darm2. The higher base
419 frequency means the number of transients in each stack is tripled and the corrected system
420 moment is nearly three times higher, resulting in more optimal performance at 90 Hz compared
421 with 30 Hz. Furthermore, the duty cycle at 30 Hz is 24 % while at 90 Hz it is 41 %. The more
422 optimal values for 90 Hz explains why the penalty for the darm2 database was only 5% while
423 for darm1 it was 15%. In other words, the information provided in the previous sections is
424 useful as it allows interpreters to assess the weaknesses of a system and to assess the quality of
425 the data generated by a system.

426 When equation (2) was applied to calculate tau from the impulse response data (e.g. Figure 9)
427 we found with the results were similar to those evident on the tau grid resulting from the
428 standard method (see Figure 8). This may be because the data was gathered with dB/dt sensors
429 and then mathematically converted into the step responses (Smith and Annan, 2000).
430 A comparison of Figures 10 and 11, shows a number of interesting features. Firstly, there are
431 less north-south asymmetry effects (highs on alternate line directions) that are evident in Figure
432 11 and secondly, the geological features running east-west and west of north west are narrower.
433 This could be a consequence of an emphasis on slightly different decay times in the figures.
434 Adjusting the window for the initial time might change the look of the image somewhat, but
435 the methods consider a different range of times and use different method so identical results
436 are not expected. If late-time information is required, then the window for V_0 can be made
437 later.

438

439 **CONCLUSIONS**

440 The practical implementation of AEM system makes tau system and processing dependent..
441 The V_{37} method proposed here is simpler and less dependent on the system and the data
442 processing . The method has some flexibility in that the initial time can be moved, to place
443 emphasis on earlier or later parts of the decay. The resulting tau grids can be used to identify
444 areas that can be used for more sophisticated processing or inversion methods.

445

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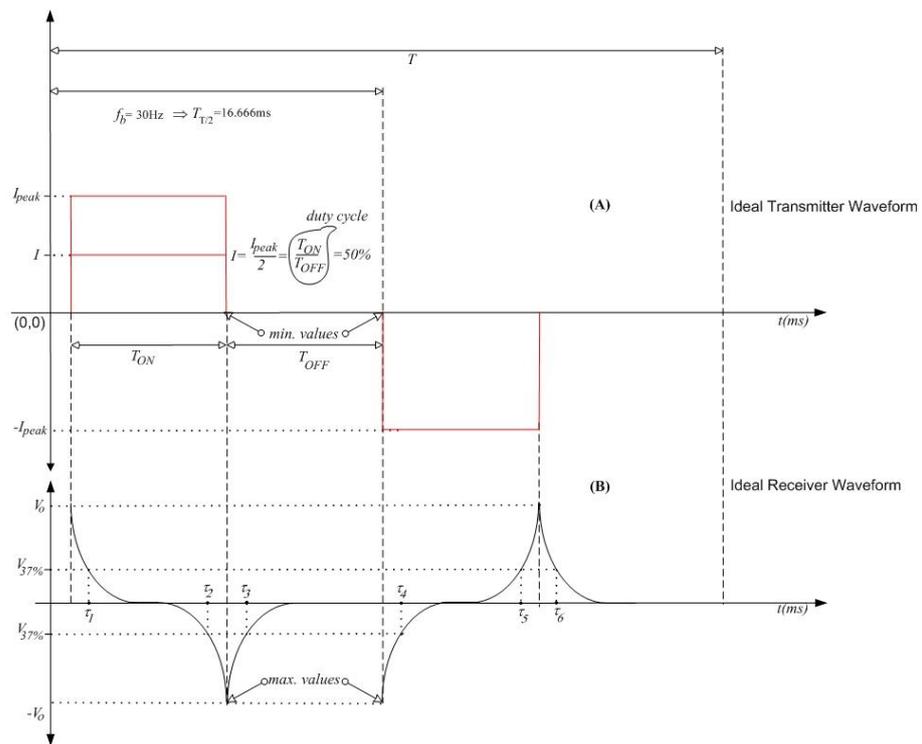
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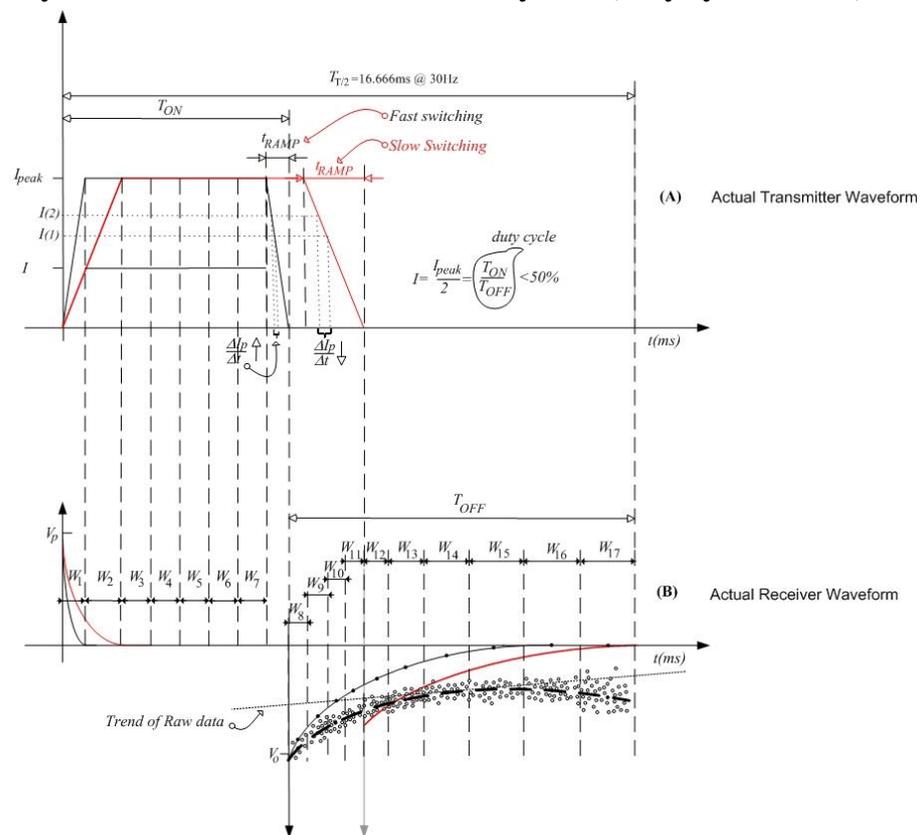
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521

522

Figure 1. Full-cycle of ideal waveforms of an EM system (duty cycle = 50%)

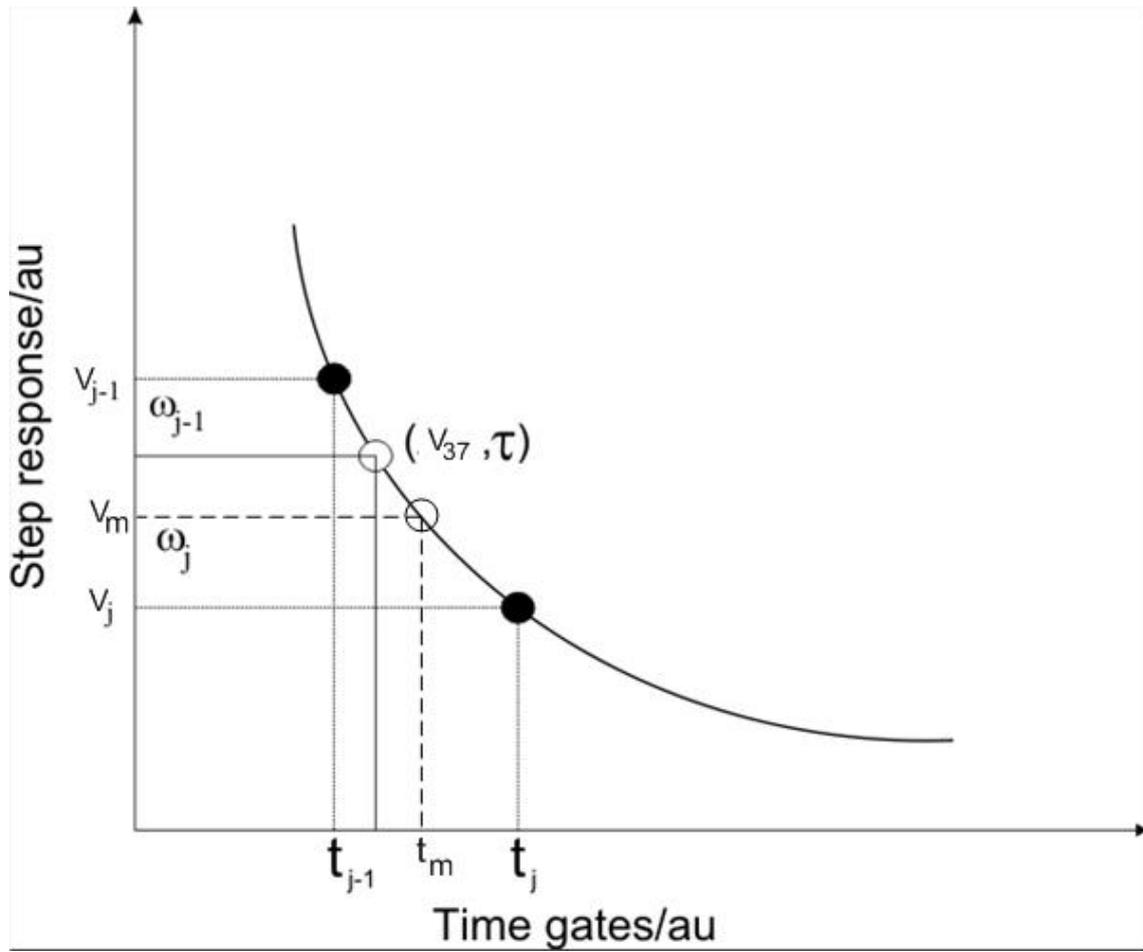


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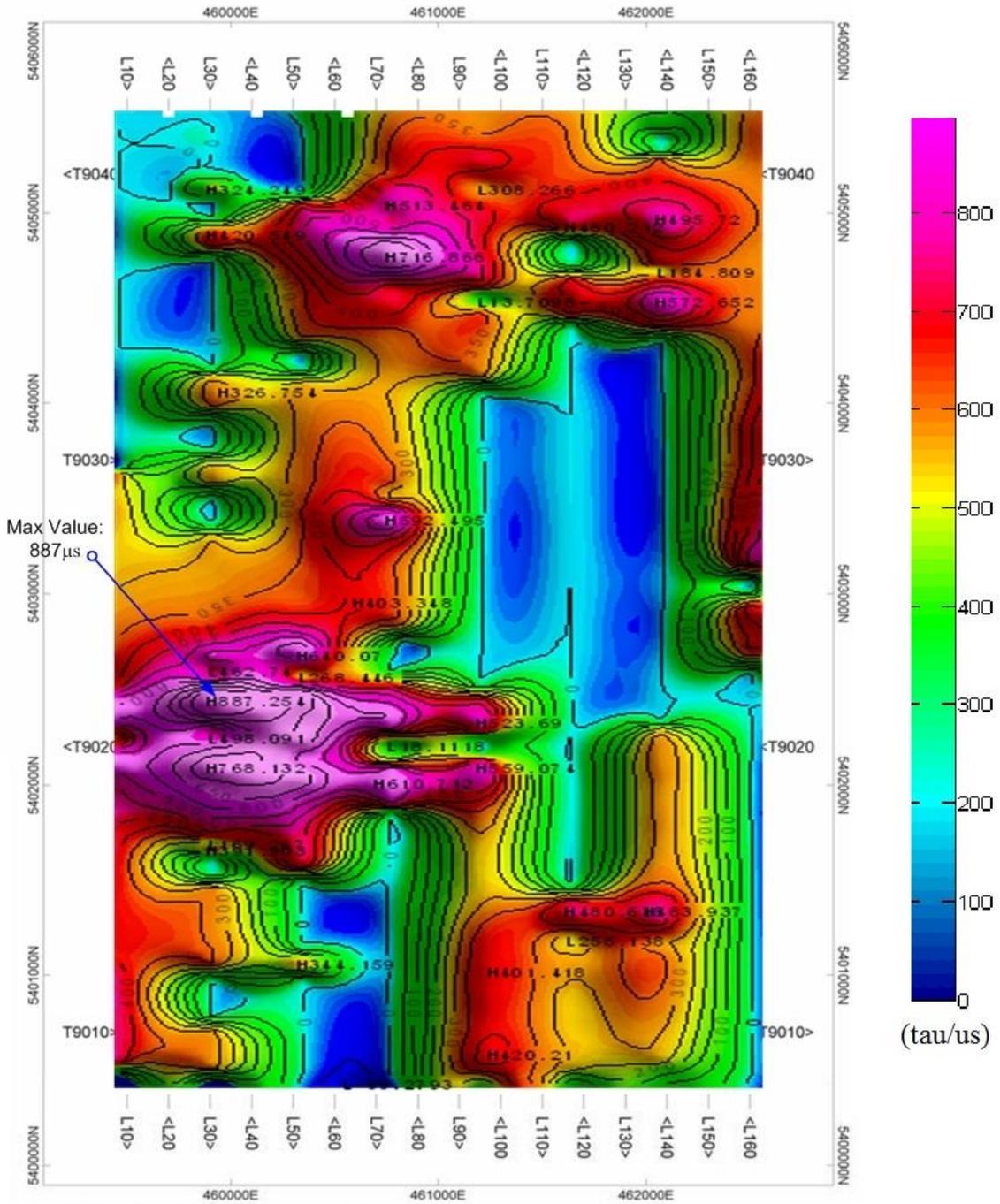
Figure 2. Half-cycle of actual waveforms of an AEM system driven by square pulse (duty cycle <50%)



526

527 **Figure 3 Schematics of V37-method to calculate τ from the time window table of AEM**
 528 **system.**

529



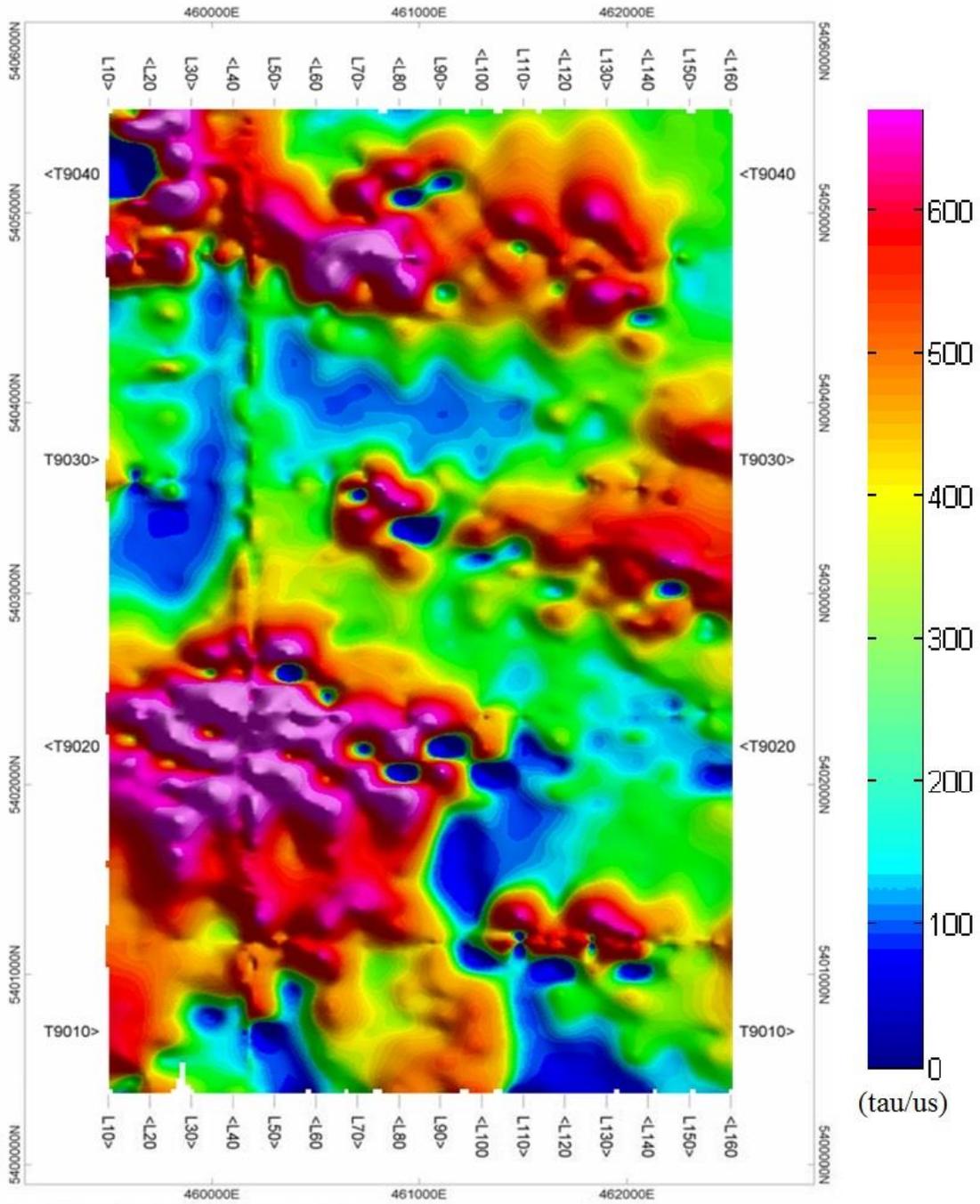
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536 **Figure 5. The tau grid from W8 using V37-method (tol.=0%) from Bz channel of darm1**

537

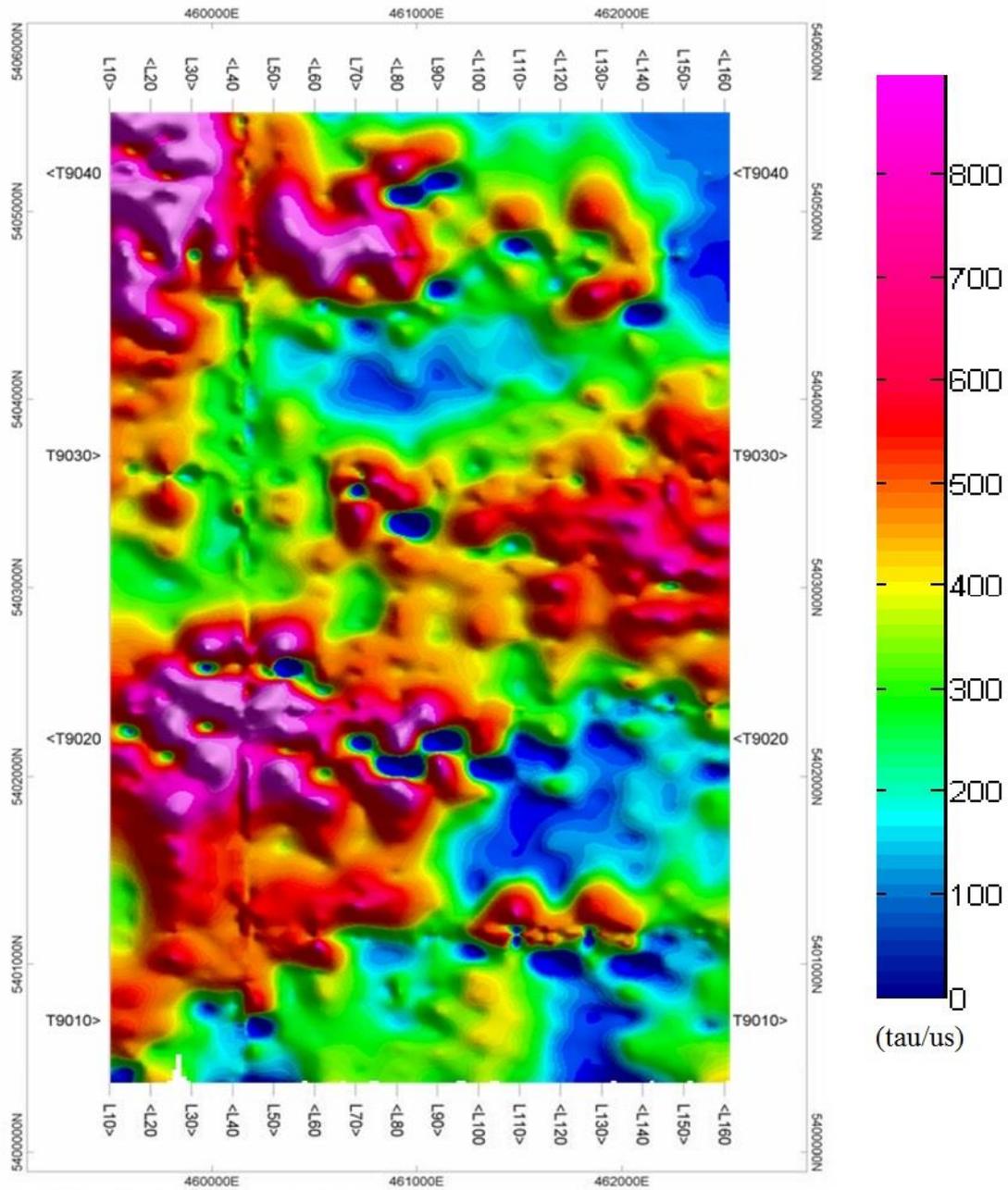
EM database [OGS (2003)]

538



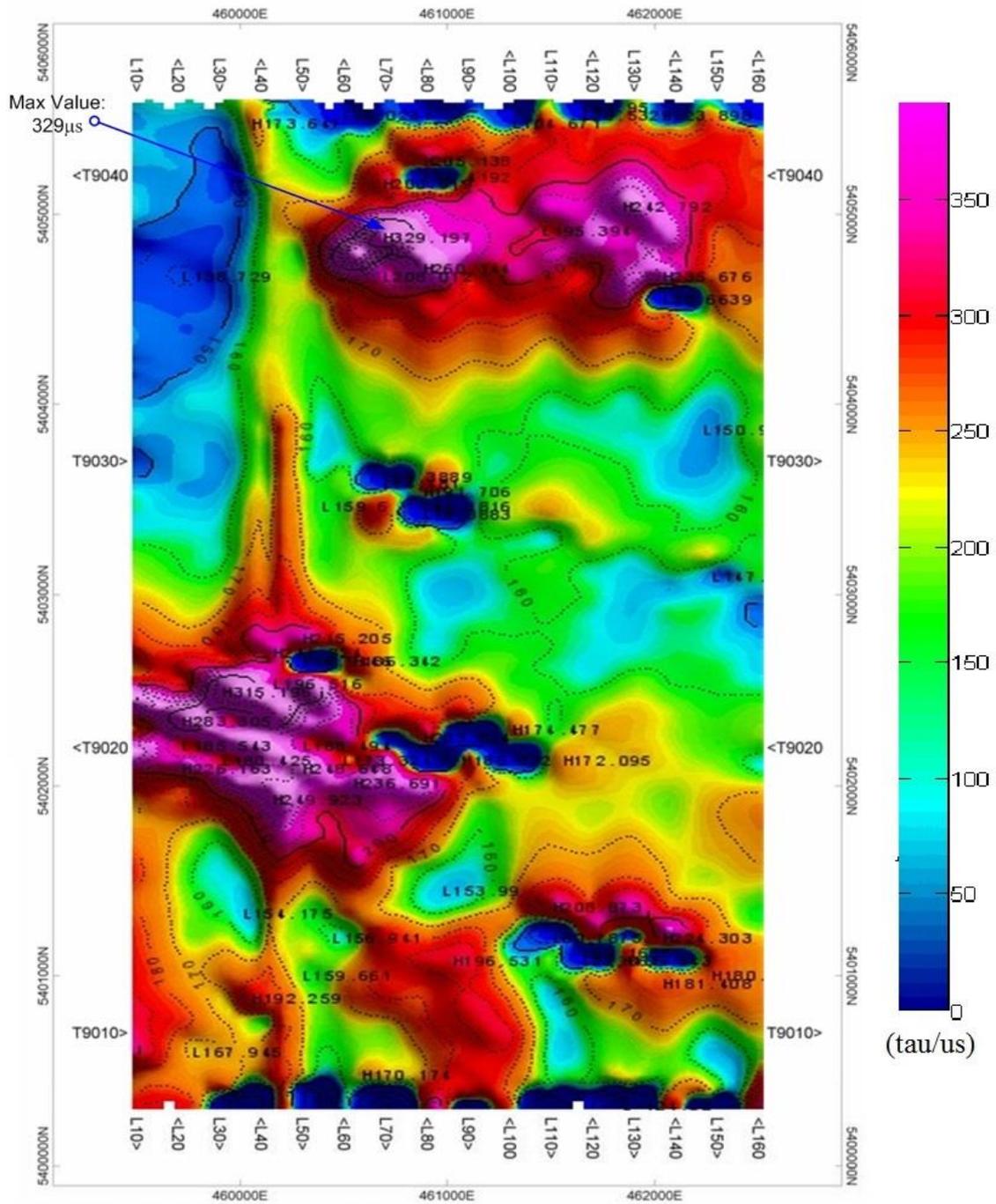
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540 **Figure 6. The tau grid from W8 using V37– method (tol.=5%) of the Bz_channel of**
541 **darm2 EM database [OGS (2004)].**



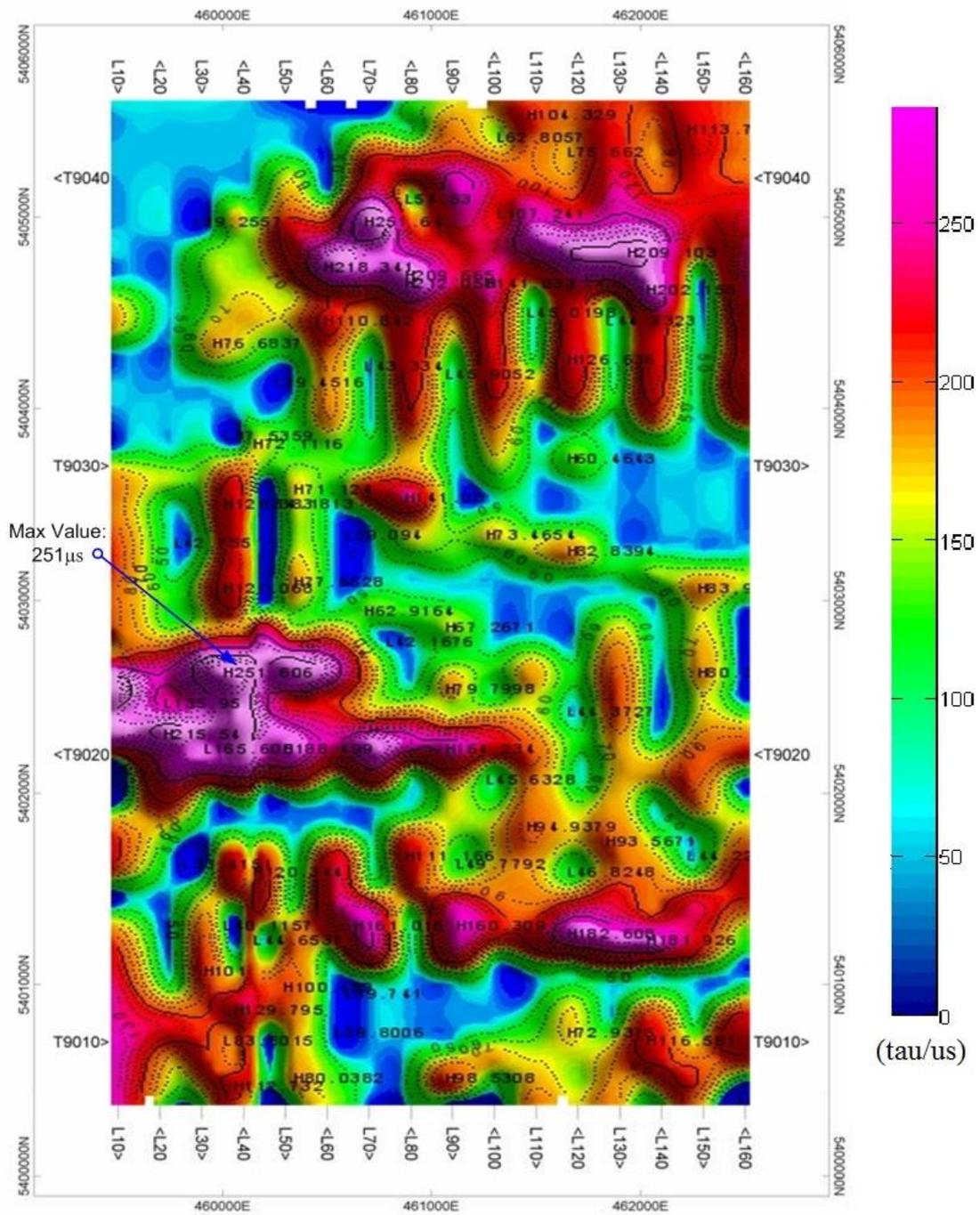
542 **Figure 7. The tau grid from W9 using V37– method (tol.=5%) of the Bz_channel of**
543 **darm2 EM database [OGS (2004)].**
544
545

547 **Figure 8. The tau grid from W8 using the standard method applied to fitting of the**
 548 **dBz/dt_channel of the darm2 EM database [OGS (2004)].**



549
 550
 551 **Figure 9. The tau grid from W8 using V37-method (tol.=5%) of dBz/dt_channel of**
 552 **darm2 EM database [OGS (2004)].**

553

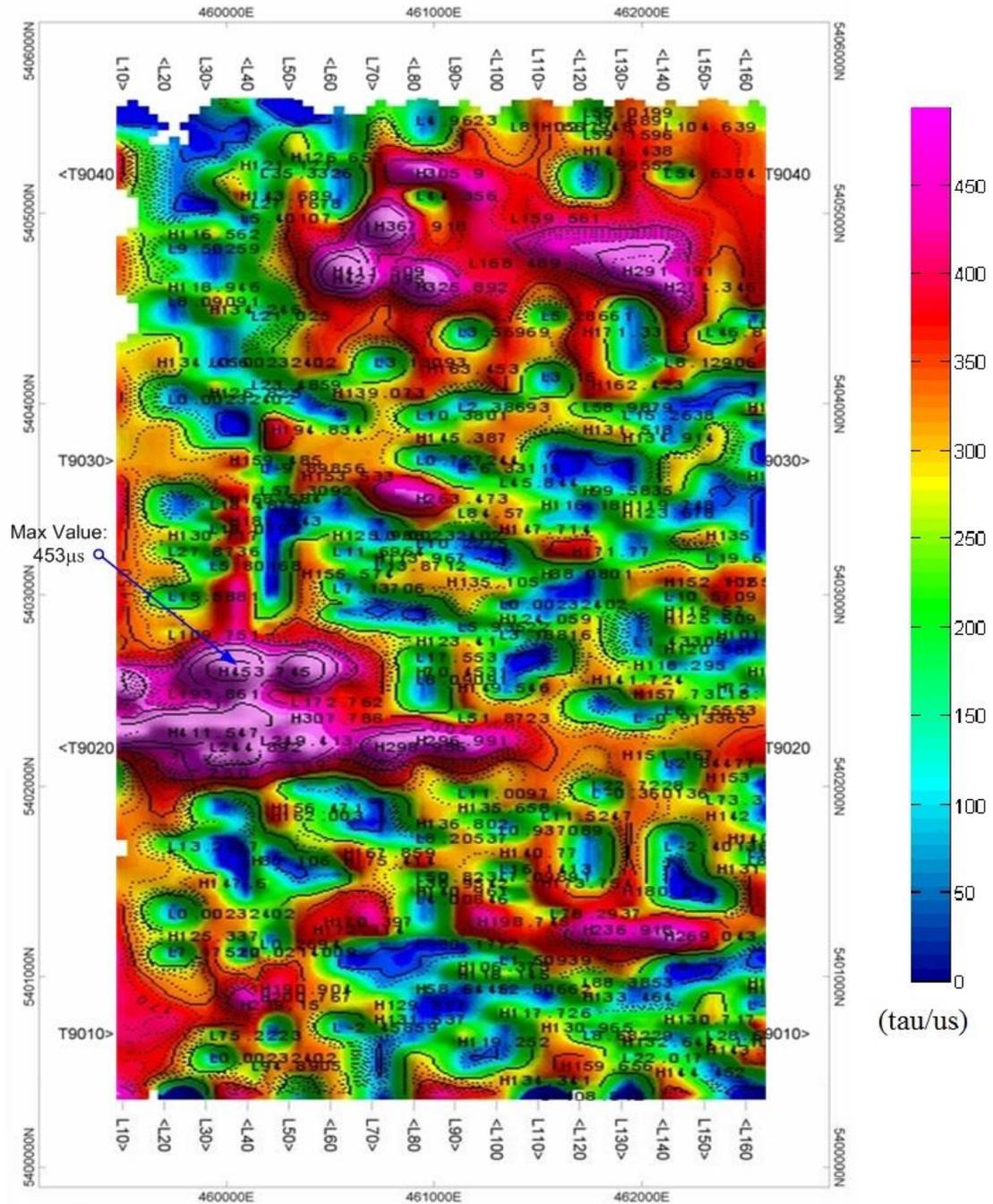


554

555

Figure 10. The tau grid from W8 using the standard method of fitting the dBx/dt_channel of the darm2 EM database [OGS (2004)].

556



557

558 **Figure 11. The tau grid from W8 using V37-method (tol.=5%) of dBx/dt_channel of**

559

darm2 EM database [OGS (2004)].

560

561 **EQUATIONS**

562

563
$$V_z(t_j) = \begin{cases} V_0 e^{-t_j/\tau} \\ dV_z(t_j)/dt = \left(\frac{V_0}{\tau} e^{-t_j/\tau} \right) \end{cases} \quad \begin{array}{l} \text{Step Response} \\ \text{Impulse Response} \end{array} \quad (1)$$

564
$$\tau = \omega_j t_j + \omega_{j-1} t_{j-1} \quad (2)$$

565
$$\omega_{j-1} = \left(1 - \frac{V_{37} - V_{j-1}}{V_{j-1} - V_j} \right) \quad (3)$$

566
$$\omega_j = \left(1 - \frac{V_j - V_{37}}{V_{j-1} - V_j} \right) \quad (4)$$

567

568 **TABLE**

569

Survey	T (s)	f_b , (Hz)	v_h (m/s)	d (m)	N_t	N_s	Sample interval (μs)	m^* Am ²
OGS2003	0.25	90	65	16	45	45	43.4	591 600
OGS2004	0.25	30	65	16	15	15	43.4	222 000

**actual moment of this survey*

570

571 **Table. Survey parameters from MEGATEM II time-domain electromagnetic system**
572 **over**

573

the Reid-Mahaffy test site